

Explicit generators in rectangular affine \mathcal{W} -algebras of type A

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Abstract

We produce in an explicit form free generators of the affine \mathcal{W} -algebra of type A associated with a nilpotent matrix whose Jordan blocks are of the same size. This includes the principal nilpotent case and we thus recover the quantum Miura transformation of Fateev and Lukyanov.

1 Main results

Let \mathfrak{g} be a reductive Lie algebra over \mathbb{C} equipped with a symmetric invariant bilinear form κ and let f be a nilpotent element of \mathfrak{g} . The corresponding *affine \mathcal{W} -algebra* $\mathcal{W}^\kappa(\mathfrak{g}, f)$ is defined by the generalized quantized Drinfeld–Sokolov reduction; see [5], [7] and [8].

In this note we take $\mathfrak{g} = \mathfrak{gl}_N$. The Jordan type of a nilpotent element $f \in \mathfrak{gl}_N$ is a partition of N . We will work with the elements f corresponding to partitions of the form (l^n) so that the associated Young diagram is the $n \times l$ rectangle with $nl = N$. Our main result is an explicit construction of free generators of the \mathcal{W} -algebra $\mathcal{W}^\kappa(\mathfrak{g}, f)$. Moreover, we calculate the images of these generators with respect to the *Miura transformation*. In particular, if f is the principal nilpotent (i.e., $n = 1$) we thus reproduce the description of the \mathcal{W} -algebra due to Fateev and Lukyanov [4]. The results can be regarded as ‘affine analogues’ of the construction of the corresponding *finite \mathcal{W} -algebras* originated in [2], [10] and extended to arbitrary nilpotent elements f in [3].

To describe the results in more detail, identify \mathfrak{g} with the tensor product of \mathfrak{gl}_l and \mathfrak{gl}_n via the isomorphism $\mathfrak{gl}_l \otimes \mathfrak{gl}_n \rightarrow \mathfrak{g}$ defined by

$$e_{ij} \otimes e_{rs} \mapsto e_{(i-1)n+r, (j-1)n+s}, \quad (1.1)$$

where the e_{ij} denote the standard basis elements of the corresponding general linear Lie algebras. Set

$$f_l = \sum_{i=1}^{l-1} e_{i+1i} \in \mathfrak{gl}_l$$

and

$$f = f_l \otimes I_n = \sum_{i=1}^{l-1} \sum_{j=1}^n e_{in+j, (i-1)n+j} \in \mathfrak{g},$$

where $I_n \in \mathfrak{gl}_n$ is the identity matrix. The matrix f is a nilpotent element of \mathfrak{g} of Jordan type (l^n) . Let

$$\mathfrak{gl}_l = \bigoplus_{p \in \mathbb{Z}} (\mathfrak{gl}_l)_p$$

be the standard principal grading of \mathfrak{gl}_l , obtained by defining the degree of e_{ij} to be equal to $j - i$. Set

$$\mathfrak{gl}_{l, \leq 0} = \bigoplus_{p \leq 0} (\mathfrak{gl}_l)_p \quad \text{and} \quad \mathfrak{gl}_{l, < 0} = \bigoplus_{p < 0} (\mathfrak{gl}_l)_p.$$

The isomorphism (1.1) then induces the \mathbb{Z} -grading on \mathfrak{g} ,

$$\mathfrak{g} = \bigoplus_{p \in \mathbb{Z}} \mathfrak{g}_p, \quad \mathfrak{g}_p = (\mathfrak{gl}_l)_p \otimes \mathfrak{gl}_n,$$

which is a *good grading* for f in the sense of [7]. We also set

$$\mathfrak{b} = \bigoplus_{p \leq 0} \mathfrak{g}_p = \mathfrak{gl}_{l, \leq 0} \otimes \mathfrak{gl}_n \quad \text{and} \quad \mathfrak{m} = \bigoplus_{p < 0} \mathfrak{g}_p = \mathfrak{gl}_{l, < 0} \otimes \mathfrak{gl}_n. \quad (1.2)$$

For any $k \in \mathbb{C}$, we let κ be any symmetric invariant bilinear form on \mathfrak{g} such that

$$\kappa(x, y) = k \operatorname{tr}(xy) \quad \text{for } x, y \in \mathfrak{sl}_N \subset \mathfrak{gl}_N. \quad (1.3)$$

For elements $x, y \in \mathfrak{b}$ set

$$\kappa_{\mathfrak{b}}(x, y) = \kappa(x, y) + \frac{1}{2} \operatorname{tr}_{\mathfrak{g}}(\operatorname{ad} x \operatorname{ad} y) - \frac{1}{2} \operatorname{tr}_{\mathfrak{g}_0} p_0(\operatorname{ad} x \operatorname{ad} y),$$

where p_0 denotes the restriction of the operator to \mathfrak{g}_0 . Then $\kappa_{\mathfrak{b}}$ defines a symmetric invariant bilinear form on \mathfrak{b} .

Example 1.1. Let

$$\kappa(x, y) = \frac{k}{2N} \operatorname{tr}_{\mathfrak{g}}(\operatorname{ad} x \operatorname{ad} y) = k \left(\operatorname{tr}(xy) - \frac{1}{N} \operatorname{tr}(x) \operatorname{tr}(y) \right), \quad x, y \in \mathfrak{g}.$$

Then for $i \geq i'$ and $j \geq j'$ we have

$$\begin{aligned} & \kappa_{\mathfrak{b}}(e_{ii'} \otimes e_{pq}, e_{jj'} \otimes e_{rs}) \\ &= \delta_{ii'} \delta_{jj'} \left((k + nl) (\delta_{ij} \delta_{ps} \delta_{qr} - \frac{1}{nl} \delta_{pq} \delta_{rs}) - n \delta_{ij} (\delta_{ps} \delta_{qr} - \frac{1}{n} \delta_{pq} \delta_{rs}) \right) \end{aligned}$$

with $N = nl$, as before. □

Let $\widehat{\mathfrak{b}} = \mathfrak{b}[t, t^{-1}] \oplus \mathbb{C}\mathbf{1}$ be the Kac–Moody affinization of \mathfrak{b} with respect to the cocycle $\kappa_{\mathfrak{b}}$, and let $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$ be the universal affine vertex algebra associated with \mathfrak{b} and $\kappa_{\mathfrak{b}}$ [6]:

$$V^{\kappa_{\mathfrak{b}}}(\mathfrak{b}) = U(\widehat{\mathfrak{b}}) \otimes_{U(\mathfrak{b}[t] \oplus \mathbb{C}\mathbf{1})} \mathbb{C},$$

where \mathbb{C} is regarded as the one-dimensional representation of $\mathfrak{b}[t] \oplus \mathbb{C}\mathbf{1}$ on which $\mathfrak{b}[t]$ acts trivially and $\mathbf{1}$ acts as 1. Note that by the Poincaré–Birkhoff–Witt theorem, $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$ is isomorphic to $U(\mathfrak{b}[t^{-1}]t^{-1})$ as a vector space.

Due to [8, 9], the \mathcal{W} -algebra $\mathcal{W}^{\kappa}(\mathfrak{g}, f)$ can be realized as a vertex subalgebra of $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$. Our aim is to give explicit description of the generators of $\mathcal{W}^{\kappa}(\mathfrak{g}, f)$ inside $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$. We will use the identification

$$\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1} \otimes \mathfrak{gl}_n \cong \mathfrak{b}[t^{-1}]t^{-1},$$

defined by

$$e_{ji}[-m] \otimes e_{pq} \mapsto e_{(j-1)n+p, (i-1)n+q}[-m], \quad m \geq 1,$$

for $1 \leq i \leq j \leq l$ and $1 \leq p, q \leq n$, where we write $x[r] = x t^r$ for any $r \in \mathbb{Z}$.

By analogy with [3, Sec. 12], consider the tensor algebra $T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1})$ of the vector space $\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}$ and let M_n denote the matrix algebra with the basis formed by the matrix units e_{ij} , $1 \leq i, j \leq n$. Define the algebra homomorphism

$$\mathcal{T} : T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}) \rightarrow M_n \otimes U(\mathfrak{b}[t^{-1}]t^{-1}), \quad x \mapsto \mathcal{T}(x) = \sum_{i,j=1}^n e_{ij} \otimes \mathcal{T}_{ij}(x)$$

by setting

$$\mathcal{T}_{ij}(x) = x \otimes e_{ji} \in \mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1} \otimes \mathfrak{gl}_n = \mathfrak{b}[t^{-1}]t^{-1}$$

for $x \in \mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}$. By definition, for any $x, y \in T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1})$ we have

$$\mathcal{T}_{ij}(xy) = \sum_{r=1}^n \mathcal{T}_{ir}(x) \mathcal{T}_{rj}(y) = \sum_{r=1}^n (x \otimes e_{ri})(y \otimes e_{jr}).$$

Let us equip the tensor product space $T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau]$ with an associative algebra structure in such a way that the natural embeddings

$$T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}) \hookrightarrow T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau] \quad \text{and} \quad \mathbb{C}[\tau] \hookrightarrow T(\mathfrak{gl}_{l, \leq 0}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau]$$

are algebra homomorphisms and the generator τ satisfies the relations

$$[\tau, x[-m]] = mx[-m-1] \quad \text{for } x \in \mathfrak{gl}_{l, \leq 0} \quad \text{and} \quad m \in \mathbb{Z}.$$

Furthermore, the tensor product space $U(\mathfrak{b}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau]$ will also be considered as an associative algebra in a similar way. We will extend \mathcal{T} to the algebra homomorphism

$$\mathcal{T} : T(\mathfrak{gl}_{\leq 0}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau] \rightarrow M_n \otimes U(\mathfrak{b}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau]$$

by setting $\mathcal{T}_{ij}(uS) = \mathcal{T}_{ij}(u)S$ for $u \in T(\mathfrak{gl}_{\leq 0}[t^{-1}]t^{-1})$ and any polynomial $S \in \mathbb{C}[\tau]$.

Set $\alpha = k + n(l - 1)$ and consider the matrix

$$B = \begin{bmatrix} \alpha\tau + e_{11}[-1] & -1 & 0 & \dots & 0 \\ e_{21}[-1] & \alpha\tau + e_{22}[-1] & -1 & \dots & 0 \\ \vdots & \vdots & \ddots & & \vdots \\ e_{l-11}[-1] & e_{l-22}[-1] & \dots & \alpha\tau + e_{l-1l-1}[-1] & -1 \\ e_{l1}[-1] & e_{l2}[-1] & \dots & \dots & \alpha\tau + e_{ll}[-1] \end{bmatrix}$$

with entries in $T(\mathfrak{gl}_{\leq 0}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau]$. Its *column-determinant* $\text{cdet } B$ is defined as the usual alternating sum of the products of the entries taken in the order determined by the column numbers of the entries.¹ So $\text{cdet } B$ is an element of $T(\mathfrak{gl}_{\leq 0}[t^{-1}]t^{-1}) \otimes \mathbb{C}[\tau]$ and we can write

$$\mathcal{T}_{ij}(\text{cdet } B) = \sum_{r=0}^l W_{ij}^{(r)}(\alpha\tau)^{l-r}$$

for certain coefficients $W_{ij}^{(r)}$ which are elements of $U(\mathfrak{b}[t^{-1}]t^{-1})$, and we can also regard them as elements of $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$. The following is our main result.

Theorem 1.2. *All coefficients $W_{ij}^{(r)}$ belong to the \mathcal{W} -algebra $\mathcal{W}^{\kappa}(\mathfrak{g}, f)$. Moreover, the \mathcal{W} -algebra $\mathcal{W}^{\kappa}(\mathfrak{g}, f) \subset V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$ is freely generated by the elements $W_{ij}^{(r)}$ with $1 \leq i, j \leq n$ and $r = 1, 2, \dots, l$.*

Set $\mathfrak{l} = (\mathfrak{gl}_{\mathfrak{l}})_0 \otimes \mathfrak{gl}_n \subset \mathfrak{gl}_N$. Then the projection $\mathfrak{b} \rightarrow \mathfrak{l}$ induces the vertex algebra homomorphism $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b}) \rightarrow V^{\kappa_{\mathfrak{b}}}(\mathfrak{l})$, which restricts to the map

$$\nu : \mathcal{W}^{\kappa}(\mathfrak{g}, f) \rightarrow V^{\kappa_{\mathfrak{b}}}(\mathfrak{l}),$$

called the (*quantum*) *Miura transformation*. This is an injective vertex algebra homomorphism. The following formula for the images of the elements $W_{ij}^{(r)}$ under the Miura transformation is an immediate consequence of Theorem 1.2.

Theorem 1.3. *We have*

$$\sum_{r=0}^l \nu(W_{ij}^{(r)})(\alpha\tau)^{l-r} = \mathcal{T}_{ij}\left((\alpha\tau + e_{11}[-1]) \dots (\alpha\tau + e_{ll}[-1])\right).$$

¹It is easy to verify that $\text{cdet } B$ coincides with the *row-determinant* of B defined in a similar way.

Note that the principal \mathcal{W} -algebra of type A corresponds to the case $n = 1$. The elements $W^{(r)}$ are defined via the expansion of $\text{cdet } B$,

$$\text{cdet } B = \sum_{r=0}^l W^{(r)}(\alpha\tau)^{l-r}.$$

By applying the Miura transformation we recover the formula of Fateev and Lukyanov [4].

Corollary 1.4. *The principal \mathcal{W} -algebra $\mathcal{W}^\kappa(\mathfrak{g}, f)$ is freely generated by the elements $W^{(1)}, \dots, W^{(l)}$. Moreover, we have*

$$\sum_{r=0}^l \nu(W^{(r)})(\alpha\tau)^{l-r} = (\alpha\tau + e_{11}[-1]) \dots (\alpha\tau + e_{ll}[-1]).$$

Example 1.5. Take $n = l = 2$ so that $N = 4$. We have

$$\text{cdet } B = (\alpha\tau)^2 + (e_{11}[-1] + e_{22}[-1])(\alpha\tau) + e_{11}[-1]e_{22}[-1] + e_{21}[-1] + \alpha e_{22}[-2].$$

Hence

$$\begin{aligned} W_{11}^{(1)} &= e_{11}[-1] + e_{33}[-1], & W_{22}^{(1)} &= e_{22}[-1] + e_{44}[-1], \\ W_{21}^{(1)} &= e_{12}[-1] + e_{34}[-1], & W_{12}^{(1)} &= e_{21}[-1] + e_{43}[-1], \\ W_{11}^{(2)} &= e_{11}[-1]e_{33}[-1] + e_{21}[-1]e_{34}[-1] + e_{31}[-1] + \alpha e_{33}[-2], \\ W_{22}^{(1)} &= e_{12}[-1]e_{43}[-1] + e_{22}[-1]e_{44}[-1] + e_{42}[-1] + \alpha e_{44}[-2], \\ W_{21}^{(1)} &= e_{12}[-1]e_{33}[-1] + e_{22}[-1]e_{34}[-1] + e_{32}[-1] + \alpha e_{34}[-2], \\ W_{12}^{(1)} &= e_{11}[-1]e_{43}[-1] + e_{21}[-1]e_{44}[-1] + e_{41}[-1] + \alpha e_{43}[-2]. \end{aligned}$$

For the images under the Miura transformation we have

$$\begin{aligned} \nu(W_{11}^{(1)}) &= e_{11}[-1] + e_{33}[-1], & \nu(W_{22}^{(1)}) &= e_{22}[-1] + e_{44}[-1], \\ \nu(W_{21}^{(1)}) &= e_{12}[-1] + e_{34}[-1], & \nu(W_{12}^{(1)}) &= e_{21}[-1] + e_{43}[-1], \\ \nu(W_{11}^{(2)}) &= e_{11}[-1]e_{33}[-1] + e_{21}[-1]e_{34}[-1] + \alpha e_{33}[-2], \\ \nu(W_{22}^{(1)}) &= e_{12}[-1]e_{43}[-1] + e_{22}[-1]e_{44}[-1] + \alpha e_{44}[-2], \\ \nu(W_{21}^{(1)}) &= e_{12}[-1]e_{33}[-1] + e_{22}[-1]e_{34}[-1] + \alpha e_{34}[-2], \\ \nu(W_{12}^{(1)}) &= e_{11}[-1]e_{43}[-1] + e_{21}[-1]e_{44}[-1] + \alpha e_{43}[-2]. \end{aligned}$$

Let the form κ_b be as in Example 1.1. The values $\kappa_b(x, y)$ are then given in the following table, where the columns and rows correspond to the x and y variables, respectively:

	e_{11}	e_{22}	e_{33}	e_{44}	e_{12}	e_{21}	e_{34}	e_{43}
e_{11}	$\frac{3k+8}{4}$	$-\frac{k}{4}$	$-\frac{k+4}{4}$	$-\frac{k+4}{4}$	0	0	0	0
e_{22}	$-\frac{k}{4}$	$\frac{3k+8}{4}$	$-\frac{k+4}{4}$	$-\frac{k+4}{4}$	0	0	0	0
e_{33}	$-\frac{k+4}{4}$	$-\frac{k+4}{4}$	$\frac{3k+8}{4}$	$-\frac{k}{4}$	0	0	0	0
e_{44}	$-\frac{k+4}{4}$	$-\frac{k+4}{4}$	$-\frac{k}{4}$	$\frac{3k+8}{4}$	0	0	0	0
e_{12}	0	0	0	0	0	$k+2$	0	0
e_{21}	0	0	0	0	$k+2$	0	0	0
e_{34}	0	0	0	0	0	0	0	$k+2$
e_{43}	0	0	0	0	0	0	$k+2$	0

These values can be used to calculate the operator product expansion formulas for the generators of $\mathcal{W}^\kappa(\mathfrak{g}, f)$. In particular, set

$$L = \frac{1}{2(k+4)} \left(-2(W_{11}^{(2)} + W_{22}^{(2)}) + W_{12}^{(1)}W_{21}^{(1)} + \frac{3}{4}(W_{11}^{(1)}W_{11}^{(1)} + W_{22}^{(1)}W_{22}^{(1)}) \right. \\ \left. - \frac{1}{2}W_{11}^{(1)}W_{22}^{(1)} - (k+2)(W_{11}^{(1)} + W_{22}^{(1)})' - (W_{11}^{(1)} - W_{22}^{(1)})' \right),$$

where the primes indicate the action of $\text{ad } \tau$ taking $e_{ij}[-1]$ to $e_{ij}[-2]$. Then L is the conformal vector of $\mathcal{W}^\kappa(\mathfrak{g}, f)$:

$$L(z)L(w) \sim -\frac{12k^2 + 41k + 32}{2(k+4)^2(z-w)^4} + \frac{2}{(z-w)^2}L(w) + \frac{1}{z-w}\partial L(w).$$

2 Proof of Theorem 1.2

Recall the notation (1.2) and let $\widehat{\mathfrak{a}} = \widehat{\mathfrak{a}}_0 \oplus \widehat{\mathfrak{a}}_1$ be the Lie superalgebra such that $\widehat{\mathfrak{a}}_0 = \widehat{\mathfrak{b}}$ and $\widehat{\mathfrak{a}}_1 = \mathfrak{m}[t, t^{-1}]$, where $\mathfrak{m}[t, t^{-1}]$ is regarded as the supercommutative Lie superalgebra, while

$$[x, y] = \text{ad } x(y) \quad \text{for } x \in \widehat{\mathfrak{a}}_0 \quad \text{and} \quad y \in \widehat{\mathfrak{a}}_1.$$

We will write $\psi_{ji}[-m] \otimes e_{pq}$ for the element

$$e_{ji}[-m] \otimes e_{pq} \in \mathfrak{gl}_{l, < 0}[t^{-1}]t^{-1} \otimes \mathfrak{gl}_n = \mathfrak{m}[t^{-1}]t^{-1}$$

with $m \geq 1$, when it is considered as an element of $\widehat{\mathfrak{a}}_1$.

Let $V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$ be the representation of $\widehat{\mathfrak{a}}$ induced from the one-dimensional representation of $(\mathfrak{b}[t] \oplus \mathbb{C}\mathbf{1}) \oplus \mathfrak{m}[t]t$ on which $\mathfrak{b}[t] \subset \widehat{\mathfrak{a}}_0$ and $\mathfrak{m}[t]t \subset \widehat{\mathfrak{a}}_1$ act trivially and $\mathbf{1}$ acts as 1. Then $V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$ is naturally a vertex algebra which contains $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$ as its vertex subalgebra. We will regard $V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$ as a (non-associative) algebra with respect to the (-1) -product

$$V^{\kappa_{\mathfrak{b}}}(\mathfrak{a}) \otimes V^{\kappa_{\mathfrak{b}}}(\mathfrak{a}) \rightarrow V^{\kappa_{\mathfrak{b}}}(\mathfrak{a}), \quad a \otimes b \mapsto a_{(-1)}b,$$

where the Fourier coefficients $a_{(n)}$ are defined in the usual way from the state-field correspondence map,

$$Y(a, z) = \sum_{n \in \mathbb{Z}} a_{(n)} z^{-n-1} \quad \text{for } a \in V^{\kappa_{\mathfrak{b}}}(\mathfrak{a}).$$

By [9] the \mathcal{W} -algebra is given by

$$\mathcal{W}^{\kappa}(\mathfrak{g}, f) = \{v \in V^{\kappa_{\mathfrak{b}}}(\mathfrak{b}) \mid Qv = 0\},$$

where $Q : V^{\kappa_{\mathfrak{b}}}(\mathfrak{a}) \rightarrow V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$ is the derivation of the non-associative algebra $V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$ defined by the following properties. First, Q commutes with the translation operator D of the vertex algebra $V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$, that is, $[Q, D] = 0$. Moreover, we have the commutation relations

$$\begin{aligned} [Q, e_{ji} \otimes e_{pq}] &= \sum_{a=i}^{j-1} \sum_{r=1}^n (e_{ai} \otimes e_{rq})(\psi_{ja} \otimes e_{pr}) - \sum_{a=i+1}^j \sum_{r=1}^n (\psi_{ai} \otimes e_{rq})(e_{ja} \otimes e_{pr}) \\ &\quad + \alpha \psi'_{ji} \otimes e_{pq} + e_{ij+1} \otimes e_{pq} - e_{i-1j} \otimes e_{pq} \end{aligned}$$

and

$$[Q, \psi_{ji} \otimes e_{pq}] = \frac{1}{2} \sum_{i < r < j, 1 \leq s \leq n} (\psi_{jr} \otimes e_{qs})(\psi_{ri} \otimes e_{ps}) - \frac{1}{2} \sum_{i < r < j, 1 \leq s \leq n} (\psi_{ri} \otimes e_{sp})(\psi_{jr} \otimes e_{qs}),$$

where we used the abbreviations

$$\begin{aligned} e_{ij} \otimes e_{pq} &= (e_{ij} \otimes e_{pq})[-1] \mathbf{1}, & \psi_{ij} \otimes e_{pq} &= (\psi_{ij} \otimes e_{pq})[-1] \mathbf{1}, \\ \psi'_{ij} \otimes e_{pq} &= D(\psi_{ij} \otimes e_{pq})[-1] \mathbf{1} = (\psi_{ij} \otimes e_{pq})[-2] \mathbf{1}, \end{aligned}$$

and set $\psi'_{ii} = 0$. Also, we used the fact that

$$\mathrm{tr}_{\mathfrak{m}} p_+(\mathrm{ad}(e_{ji} \otimes e_{pq}) \mathrm{ad}(e_{ij} \otimes e_{pq})) = n(l + i - j - 1)$$

for $1 \leq i < j \leq l$ and $1 \leq p, q \leq n$, where p_+ denotes the restriction of the operator to \mathfrak{m} .

Our goal now is to reduce the calculations to the principal nilpotent case. To this end, when $n = 1$, we will write $\bar{\mathfrak{a}}$ and $\bar{\mathfrak{b}}$ respectively, instead of \mathfrak{a} and \mathfrak{b} , and replace k with $k + (n - 1)(l - 1)$ in (1.3). Consequently, $V^{\kappa_{\bar{\mathfrak{b}}}}(\bar{\mathfrak{a}})$ will denote the vertex algebra $V^{\kappa_{\mathfrak{b}}}(\mathfrak{a})$ with $n = 1$ (and k replaced by $k + (n - 1)(l - 1)$). We let \bar{Q} denote the operator Q for $V^{\kappa_{\bar{\mathfrak{b}}}}(\bar{\mathfrak{a}})$. We have

$$\begin{aligned} [\bar{Q}, e_{ji}] &= \sum_{a=i}^{j-1} e_{ai} \psi_{ja} - \sum_{a=i+1}^j \psi_{ai} e_{ja} + \alpha \psi'_{ji} + \psi_{j+1i} - \psi_{ji-1}, \\ [\bar{Q}, \psi_{ji}] &= \frac{1}{2} \sum_{i < r < j} (\psi_{jr} \psi_{ri} - \psi_{ri} \psi_{jr}), \end{aligned}$$

where we used the notation $e_{ji} = e_{ji}[-1]$, $\psi_{ij} = \psi_{ij}[-1]$, $\psi'_{ij} = \psi_{ij}[-2]$, and we set $\psi'_{ii} = 0$.

We will regard $V^{\kappa_b}(\mathbf{a}) \otimes \mathbb{C}[\tau]$ as a non-associative algebra with the natural subalgebras $V^{\kappa_b}(\mathbf{a})$ and $\mathbb{C}[\tau]$ together with the relation $[\tau, u] = Du$ for $u \in V^{\kappa_b}(\mathbf{a})$. Similarly, the tensor product $V^{\kappa_{\bar{b}}}(\bar{\mathbf{a}}) \otimes \mathbb{C}[\tau]$ will be regarded as a non-associative algebra with the relation $[\tau, u] = \bar{D}u$ for $u \in V^{\kappa_{\bar{b}}}(\bar{\mathbf{a}})$, where \bar{D} denotes the translation operator of the vertex algebra $V^{\kappa_{\bar{b}}}(\bar{\mathbf{a}})$. Define the non-associative algebra homomorphism

$$\tilde{\mathcal{T}} : V^{\kappa_{\bar{b}}}(\bar{\mathbf{a}}) \otimes \mathbb{C}[\tau] \rightarrow M_n \otimes V^{\kappa_b}(\mathbf{a}) \otimes \mathbb{C}[\tau], \quad x \mapsto \tilde{\mathcal{T}}(x) = \sum_{p,q=1}^n e_{pq} \otimes \tilde{\mathcal{T}}_{pq}(x)$$

by

$$\tilde{\mathcal{T}}_{pq}(e_{ji}[-m]) = e_{ji}[-m] \otimes e_{qp}, \quad \tilde{\mathcal{T}}_{pq}(\psi_{ji}[-m]) = \psi_{ji}[-m] \otimes e_{qp} \quad \text{and} \quad \tilde{\mathcal{T}}_{pq}(\tau) = \tau.$$

We extend the definition of the column-determinant to matrices $A = [a_{ij}]$ with entries in a non-associative algebra by using right-normalized products,

$$\widetilde{\text{cdet}} A = \sum_{\sigma \in \mathfrak{S}_l} \text{sgn } \sigma \cdot a_{\sigma(1)1} (a_{\sigma(2)2} (a_{\sigma(3)3} (\dots (a_{\sigma(l-1)l-1} a_{\sigma(l)l}))). \quad (2.4)$$

Note the relation

$$\tilde{\mathcal{T}}(\widetilde{\text{cdet}} B) = \mathcal{T}(\text{cdet } B),$$

where $\widetilde{\text{cdet}} B$ is regarded as an element of $V^{\kappa_{\bar{b}}}(\bar{\mathbf{a}}) \otimes \mathbb{C}[\tau]$. The first part of Theorem 1.2 will now be implied by the following two propositions.

Proposition 2.1. *For any $a \in V^{\kappa_{\bar{b}}}(\bar{\mathbf{a}}) \otimes \mathbb{C}[\tau]$ and $1 \leq p, q \leq n$ we have the relations*

$$[Q, \tilde{\mathcal{T}}_{pq}(a)] = \tilde{\mathcal{T}}_{pq}([\bar{Q}, a]).$$

Proof. This follows immediately from the definitions of the operators Q and \bar{Q} . \square

Proposition 2.2. *We have the relation*

$$[\bar{Q}, \widetilde{\text{cdet}} B] = 0.$$

Proof. We use induction on l . For any $0 \leq s \leq l$ consider the submatrix $B^{(s)}$ of B corresponding to its last s rows and columns, which is given by

$$\begin{bmatrix} \alpha\tau + e_{l-s+1l-s+1}[-1] & -1 & 0 & \dots & 0 \\ e_{l-s+2l-s+1}[-1] & \alpha\tau + e_{l-s+2l-s+2}[-1] & -1 & \dots & 0 \\ \vdots & \vdots & \ddots & & \vdots \\ e_{l-1l-s+1}[-1] & e_{l-1l-s+2}[-1] & \dots & \alpha\tau + e_{l-1l-1}[-1] & -1 \\ e_{ll-s+1}[-1] & e_{ll-s+2}[-1] & \dots & \dots & \alpha\tau + e_{ll}[-1] \end{bmatrix}.$$

Using the definition (2.4), set $D^{(s)} = \widetilde{\text{cdet}} B^{(s)}$.

Lemma 2.3. *We have the column expansion formula*

$$D^{(s)} = \sum_{i=1}^s B_{i1}^{(s)} D^{(s-i)},$$

where $B_{ij}^{(s)}$ denotes the (i, j) entry of $B^{(s)}$. □

Suppose that $s < l$. By the induction hypothesis, the commutator $[\overline{Q}, D^{(s)}]$ equals

$$\sum_{i=1}^s \widetilde{\text{cdet}} \begin{bmatrix} 0 & -1 & 0 & \dots & 0 \\ 0 & \alpha \tau + e_{l-s+2l-s+2}[-1] & -1 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ -\psi_{l-i+1l-s}[-1] & \dots & \dots & \dots & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & e_{l-1l-s+2}[-1] & \dots & \alpha \tau + e_{l-1l-1}[-1] & -1 \\ 0 & e_{ll-s+2}[-1] & \dots & \dots & \alpha \tau + e_{ll}[-1] \end{bmatrix}$$

so that

$$[\overline{Q}, D^{(s)}] = - \sum_{i=1}^s \psi_{l-i+1l-s}[-1] D^{(i-1)}. \quad (2.5)$$

Hence, by Lemma 2.3, we have

$$[\overline{Q}, D^{(l)}] = \sum_{i=1}^l [\overline{Q}, B_{i1}^{(l)}] D^{(l-i)} + \sum_{i=1}^l B_{i1}^{(l)} [\overline{Q}, D^{(l-i)}].$$

Now we use the definition of \overline{Q} and relation (2.5) to write this expression as

$$\sum_{i=1}^l \left(\sum_{a=1}^{i-1} e_{a1} \psi_{ia} - \sum_{a=2}^i \psi_{a1} e_{ia} + \alpha \psi'_{i1} + \psi_{i+11} \right) D^{(l-i)} - \sum_{i=1}^l B_{i1}^{(l)} \sum_{a=1}^{l-i} (\psi_{l-a+1i} D^{(a-1)}),$$

where, as before, we write $e_{ij} = e_{ij}[-1]$, $\psi_{ij} = \psi_{ij}[-1]$ and $\psi'_{ij} = \psi_{ij}[-2]$ for brevity. Thus,

$$\begin{aligned} [\overline{Q}, D^{(l)}] &= \sum_{i=1}^l \sum_{a=1}^{i-1} e_{a1} (\psi_{ia} D^{(l-i)}) - \sum_{i=1}^l \sum_{a=2}^i \psi_{a1} (e_{ia} D^{(l-i)}) + \sum_{i=1}^l (\alpha \psi'_{i1} + \psi_{i+11}) D^{(l-i)} \\ &\quad - \sum_{i=1}^l B_{i1}^{(l)} \sum_{a=1}^{l-i} (\psi_{l-a+1i} D^{(a-1)}) \end{aligned}$$

which equals

$$\begin{aligned}
& -\alpha \sum_{a=2}^l \psi_{a1} \tau D^{(l-a)} - \sum_{a=2}^l \sum_{i=a}^l \psi_{a,1} (e_{ia} D^{(l-i)}) + \sum_{i=1}^l \psi_{i+1,1} D^{(l-i)} \\
& = -\sum_{a=1}^l \psi_{a1} \left(\sum_{a=i}^l (\delta_{ia} \alpha \tau + e_{ia}) D^{(l-i)} - D^{(l-a+1)} \right) \\
& = \sum_{a=1}^l \left(D^{(l-a+1)} - \sum_{i=1}^{l-a+1} B_{ia}^{(l-a+1)} D^{(l-a+1-i)} \right) = 0,
\end{aligned}$$

where the last equality holds by Lemma 2.3. Here we used the relations

$$\begin{aligned}
(e_{ji}[-m] \psi_{pq}[-n]) u &= e_{ji}[-m] (\psi_{pq}[-n] u), \\
(\psi_{pq}[-n] e_{ji}[-m]) u &= \psi_{pq}[-n] (e_{ji}[-m] u),
\end{aligned}$$

which hold under the assumption

$$u \in \text{span of } \{e_{i'j'}[-m'], \psi_{p'q'}[-n'] \mid j' > j \text{ and } q' > p\}.$$

This completes the proof of the proposition. \square

To see the second part of Theorem 1.2, consider the grading of $V^{\kappa_{\mathfrak{b}}}(\mathfrak{b})$ induced by the grading of \mathfrak{b} . One has

$$W_{ij}^{(r)} = \mathcal{T}_{ij} \left(\sum_{s=1}^{l-r+1} e_{r+s-1s}[-1] \right) + (\text{terms of higher degree}).$$

Now the elements $\sum_{s=1}^{l-r+1} e_{r+s-1s}$ with $r = 1, \dots, l$ form a basis of \mathfrak{g}_l^f and the elements

$$\sum_{s=1}^{l-r+1} e_{r+s-1s} \otimes e_{ji}, \quad r = 1, \dots, l \text{ and } i, j = 1, \dots, n,$$

form a basis of \mathfrak{g}^f . Hence the claim follows from [8, Theorem 4.1] (cf. [1, Theorem 5.5.1]) thus completing the argument.

References

- [1] T. Arakawa, *Representation theory of superconformal algebras and the Kac–Roan–Wakimoto conjecture*, Duke Math. J. **130** (2005), 435–478.
- [2] C. Briot and E. Ragoucy, *RTT presentation of finite \mathcal{W} -algebras*, J. Phys. A **34** (2001), 7287–7310.

- [3] J. Brundan and A. Kleshchev, *Shifted Yangians and finite W -algebras*, Adv. Math. **200** (2006), 136–195.
- [4] V. A. Fateev and S. L. Lykhanov, *The models of two-dimensional conformal quantum field theory with Z_n symmetry*, Internat. J. Modern Phys. A **3** (1988), 507–520.
- [5] B. Feigin and E. Frenkel, *Quantization of the Drinfeld–Sokolov reduction*, Phys. Lett. B **246** (1990), 75–81.
- [6] V. Kac, *Vertex algebras for beginners*, University Lecture Series, 10. American Mathematical Society, Providence, RI, 1997.
- [7] V. Kac, Shi-Shyr Roan and M. Wakimoto, *Quantum reduction for affine superalgebras*, Comm. Math. Phys. **241** (2003), 307–342.
- [8] V. Kac and M. Wakimoto, *Quantum reduction and representation theory of superconformal algebras*, Adv. Math. **185** (2004), 400–458.
- [9] V. Kac and M. Wakimoto, *Corrigendum to: “Quantum reduction and representation theory of superconformal algebras” [Adv. Math. 185 (2004), 400–458]*, Adv. Math. **193** (2005), 453–455.
- [10] E. Ragoucy and P. Sorba, *Yangian realisations from finite \mathcal{W} -algebras*, Comm. Math. Phys. **203** (1999), 551–572.

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